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Nonlocality in peridynamics

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Outline

- Nonlocality
 - It's not as weird as everybody thinks
- Peridynamics background
 - All-in on nonlocality
- Can nonlocality be derived or observed?
 - Long-range forces
 - Smoothed degrees of freedom (homogenization)
 - Multiple pathways for flux
 - Wave dispersion



Do we ask too much of the local theory of continuum mechanics?

What peridynamics seeks to accomplish



- Treat material points on or off of evolving discontinuities with the same equations.
- Include long-range forces in the basic equations.
- Fit all this into a thermodynamic framework that's consistent with the mechanics.



Peridynamic simulation



Metallic glass crack tip*

*Hofmann et al, Nature (2008)

Nonlocality: Not as weird as everybody thinks



Discretized numerical methods are nonlocal



- Node *i* interacts directly with node *j* through the finite element equations.
- Interaction is across a finite distance Δx .
- This is a form of nonlocality.
 - Notwithstanding that the result converges to the local result as $\Delta x \rightarrow 0$.



Local PDEs get themselves into trouble

• Classical (Cauchy) PDE:

$$ho \ddot{\mathbf{u}} =
abla \cdot oldsymbol{\sigma} \left(rac{\partial \mathbf{u}}{\partial \mathbf{x}}
ight) + \mathbf{b} \cdot oldsymbol{\sigma}$$

- Many material models $\sigma(\cdot)$ evolve into deformations that are incompatible with the fundamental assumptions.
 - Phase boundaries, shock waves, cracks, ...
- Can't directly treat some important physical effects.
 - Wave dispersion, surface energy, microstructure evolution, long-range forces, . . .
- People often take drastic measures if they want to work with this PDE.
 - Element deletion, ...



Nonlocality: Not as weird as everybody thinks

These drastic measures often involve nonlocality

• Example: Artificial viscosity spreads out a shock wave and dissipates energy.

$$ho \ddot{\mathbf{u}} =
abla \cdot oldsymbol{\sigma} \left(rac{\partial \mathbf{u}}{\partial \mathbf{x}}
ight) + \gamma (
abla \cdot \dot{\mathbf{u}})^2 + \mathbf{b}.$$

- It avoids the need to apply jump conditions across an ideal shock.
- It allows converntional discretization to be used "within" a shock.
- By spreading out a shock it introduces a length scale.
- This is a type of nonlocality.





• J. Von Neumann & R. D. Richtmyer, J. Appl. Phys. 21 (1950). 232

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Peridynamics goes all-in on nonlocality

Classification of some theories with respect to local/nonlocality:



- Every fundamental relation in peridynamics is nonlocal in space:
 - Transport
 - Conservation
 - Material models

Peridynamic* momentum balance

- Any point x interacts directly with other points within a distance δ called the "horizon."
- The material within a distance δ of x is called the "family" of x, \mathcal{H}_x .





- If **f** satisfies f(x, q) = -f(q, x) for all x, q then linear momentum is conserved.
- SS, JMPS (2000)
- * Peri (near) + dyne (force)





Formalism for nonlocal interactions: States

• A state is a mapping whose domain is all the bonds ξ in a family.

$${f A}\langle {m \xi}
angle = {f something} \qquad orall {m \xi}\in {\cal H}.$$



• Deformation state...

 $\underline{\mathbf{Y}}[\mathbf{x}]\langle \mathbf{q}-\mathbf{x}\rangle=\mathbf{y}(\mathbf{q})-\mathbf{y}(\mathbf{x})=\text{deformed image of the bond}$

States: Nonlocal analogues of second order tensors

- Classical theory uses tensors (linear mappings from vectors to vectors).
- Peridynamics uses states (nonlinear mappings from vectors to vectors).





Peridynamic vs. local equations

• Structurally similar but with states instead of local operators.

Relation	Peridynamic theory	Standard theory
Kinematics	$\underline{\mathbf{Y}}\langle \mathbf{q}-\mathbf{x} angle = \mathbf{y}(\mathbf{q})-\mathbf{y}(\mathbf{x})$	$\mathbf{F}(\mathbf{x}) = rac{\partial \mathbf{y}}{\partial \mathbf{x}}(\mathbf{x})$
Linear momentum balance	$\rho \ddot{\mathbf{y}}(\mathbf{x}) = \int_{\mathcal{H}} \left(\mathbf{t}(\mathbf{q}, \mathbf{x}) - \mathbf{t}(\mathbf{x}, \mathbf{q}) \right) dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x})$	$ ho \ddot{\mathbf{y}}(\mathbf{x}) = \nabla \cdot \boldsymbol{\sigma}(\mathbf{x}) + \mathbf{b}(\mathbf{x})$
Constitutive model	$\mathbf{t}(\mathbf{q},\mathbf{x}) = \underline{\mathbf{T}}\langle \mathbf{q} - \mathbf{x} \rangle, \qquad \underline{\mathbf{T}} = \underline{\hat{\mathbf{T}}}(\underline{\mathbf{Y}})$	$oldsymbol{\sigma} = \hat{oldsymbol{\sigma}}(\mathbf{F})$
Angular momentum balance	$\int_{\mathcal{H}} \underline{\mathbf{Y}} \langle \mathbf{q} - \mathbf{x} \rangle \times \underline{\mathbf{T}} \langle \mathbf{q} - \mathbf{x} \rangle \ dV_{\mathbf{q}} = 0$	$oldsymbol{\sigma} = oldsymbol{\sigma}^T$
Elasticity	$\mathbf{\underline{T}}=W_{\mathbf{\underline{Y}}}$ (Fréchet derivative)	$oldsymbol{\sigma} = W_{\mathbf{F}}$ (tensor gradient)
First law	$\dot{\varepsilon} = \underline{\mathbf{T}} \bullet \underline{\dot{\mathbf{Y}}} + q + r$	$\dot{\varepsilon} = \boldsymbol{\sigma} \cdot \dot{\mathbf{F}} + q + r$
$\mathbf{\underline{T}} ullet \dot{\mathbf{Y}} := \int_{\mathcal{H}} \mathbf{\underline{T}} \langle oldsymbol{\xi} angle \ dV_{oldsymbol{\xi}}$		



Damage

- Damage is usually treated through *bond breakage*.
- After a bond ξ breaks according to some criterion, it no longer carries any force.
- Typical breakage criterion: prescribed *critical bond strain* s₀:

$$s = rac{|\mathbf{Y}\langle \boldsymbol{\xi}
angle| - |\boldsymbol{\xi}|}{|\boldsymbol{\xi}|}$$
 bond strain.

 $s >= s_0$ at some time t_0

means the bond remains broken for all $t \ge t_0$.



Autonomous crack growth



• Cracks do what they want (grow, arrest, branch, curve, oscillate, ...)



• SS & Askari, Computers and Structures (2005)